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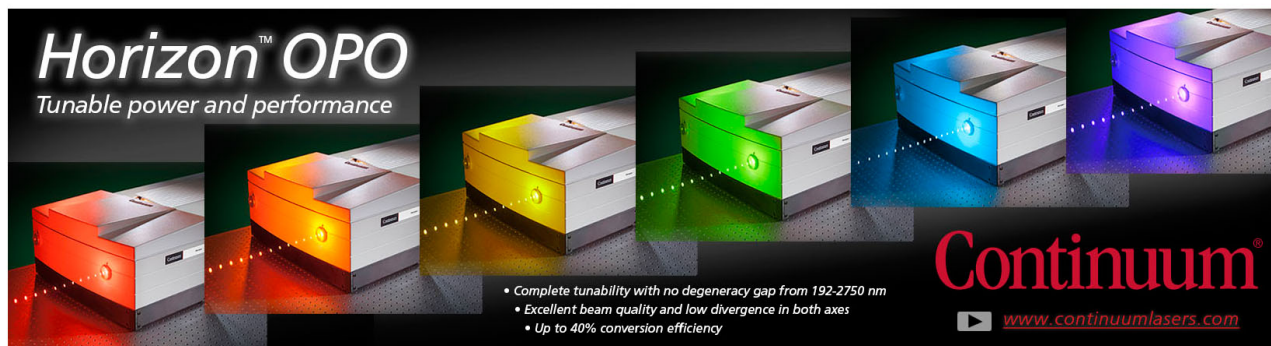
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# A comprehensive scenario of the thermodynamic anomalies of water using the TIP4P/2005 model

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The striking behavior of water has deserved it to be referred to as an “anomalous” liquid. The water anomalies are greatly amplified in metastable (supercooled and/or stretched) regions. This makes difficult a complete experimental description since, beyond certain limits, the metastable phase necessarily transforms into the stable one. Theoretical interpretation of the water anomalies could then be based on simulation results of well validated water models. But the analysis of the simulations has not yet reached a consensus. In particular, one of the most popular theoretical scenarios—involving the existence of a liquid-liquid critical point (LLCP)—is disputed by several authors. In this work, we propose to use a number of exact thermodynamic relations which may shed light on this issue. Interestingly, these relations may be tested in a region of the phase diagram which is *outside* the LLCP thus avoiding the problems associated to the coexistence region. The central property connected to other water anomalies is the locus of temperatures at which the density along isobars attain a maximum (TMD line) or a minimum (TmD). We have performed computer simulations to evaluate the TMD and TmD for a successful water model, namely, TIP4P/2005. We have also evaluated the vapor-liquid (VL) spinodal in the region of large negative pressures. The shape of these curves and their connection to the extrema of some response functions, in particular the isothermal compressibility and heat capacity at constant pressure, provides very useful information which may help to elucidate the validity of the theoretical proposals. In this way, we are able to present for the first time a comprehensive scenario of the thermodynamic water anomalies for TIP4P/2005 and their relation to the vapor-liquid spinodal. The overall picture shows a remarkable similarity with the corresponding one for the ST2 water model, for which the existence of a LLCP has been demonstrated in recent years. It also provides a hint as to where the long-sought for extrema in response functions might become accessible to experiments. *Published by AIP Publishing.* [<http://dx.doi.org/10.1063/1.4960185>]

## I. INTRODUCTION

The physical properties of water at ambient conditions are markedly different from those of other liquids. It is widely known that the density of liquid water at a fixed pressure exhibits a maximum at the so-called temperature of maximum density (TMD). In particular, at atmospheric pressure, the TMD is approximately 4 °C. Below this temperature, the water's expansivity,  $\alpha$ , is negative in striking contrast with “normal” liquids where  $\alpha$  is always positive. Other thermodynamic response functions such as the isothermal compressibility,  $\kappa_T$ , or the isobaric heat capacity  $C_p$  also show an unusual behavior.<sup>1</sup> On supercooling, these anomalies are enhanced. In particular, the isothermal compressibility seems to diverge at 228 K.<sup>2</sup>

Several scenarios have been proposed to account for the water anomalies and their magnification at temperatures below the melting point.<sup>3–6</sup> In the stability limit conjecture (SLC),<sup>3</sup> the increase of the response functions in the supercooled region

is ascribed to a continuous retracing line of instability that delimits the supercooled and stretched metastable states. The second critical point scenario assumes that there is a liquid-liquid coexistence terminating at a critical point (LLCP).<sup>4</sup> The response functions would reach a peak and converge towards the Widom line (the line of the maxima of the correlation length) emanating from the LLCP.<sup>7,8</sup> Finally, it has been shown that thermodynamic consistency explains the existence of peaks in the response functions as a mere consequence of the presence of density anomalies. In this case, there is no singular behavior, hence the name singularity-free (SF) scenario.<sup>5</sup> Notice that the SF interpretation can also be seen as the second critical-point hypothesis with the LLCP occurring at zero temperature.

The experimental testing of these scenarios is extremely difficult, if not impossible, because of the difficulties to access the “no man's land” region that lies below the temperature at which water spontaneously freezes. Crystallization may be inhibited by confining water in nanosized samples, but it is

unclear whether surface effects could influence the outcome of the experiments.<sup>9–11</sup> Different experimental approaches have been proposed to circumvent the problems associated with the bulk water no man's land<sup>12–22</sup> (see also a recent review<sup>23</sup> on this topic). They have provided very important information on the behavior of water at extreme conditions (deep supercooling and/or high negative pressures). Although the experiments seem to be consistent and support the appearance of a liquid-liquid transition, their interpretation is not conclusive.

Theoretical work has shown that the slope of the TMD loci determines the behavior of the thermodynamic response functions. Exact thermodynamic relations relate the shape of the TMD to that of other water anomalies. It is well known that the TMD of liquid water is a negatively sloped function in a  $p$ - $T$  diagram. If the negative slope would extend to large negative pressures (SLC scenario), it would meet the vapor-liquid (VL) spinodal. In such a case, thermodynamic consistency requires<sup>24</sup> that the spinodal would retrace at the intersection point. On the other hand, the TMD could change its slope leading to a nose-shaped function (LLCP and SF scenarios). It has been demonstrated that a positively sloped TMD line cannot cross a positively sloped spinodal in a thermodynamically consistent phase diagram and that the turning point of the TMD must intersect the locus of isothermal compressibility extrema.<sup>5</sup> In summary, the study of the TMD of stretched water may give insight to the validity of the hypothesis proposed to explain the water anomalies. Recent measurements of the speed of sound<sup>20</sup> have enabled to extend considerably our knowledge of the TMD in the region of large negative pressures.<sup>25</sup> These results indicate that the slope of the TMD becomes increasingly more negative as the pressure decreases and strongly suggest that the experimental TMD is about to reach a retracing point. Unfortunately, bubble nucleation prevents carrying this study further.

Given the experimental difficulties, it is clear that molecular simulation may be an alternative for our purpose. Most of the computer simulation studies using realistic water models seem to support the existence of the liquid-liquid separation. The seminal work of Poole *et al.*<sup>4</sup> focused on the ST2 water model.<sup>26</sup> Some of the features of the ST2 model allow a thorough investigation of the supercooled region. Because of this, the model has been widely used in the study of the liquid-liquid phase transition. Most simulations using ST2<sup>7,8,27–37</sup> seem to have unambiguously demonstrated the existence of a LLCP although this interpretation has been challenged by Limmer and Chandler.<sup>38,39</sup> But the advantages of the model for the study of the supercooled region (among them, a high value for the TMD) are closely related to its major drawback: ST2 is known to produce an over-structured liquid compared to real water. Thus, there is no compelling evidence that the behavior of metastable water can be described by ST2.<sup>29</sup>

Alternative successful water models can indeed be found in the literature though they are not free of objections. SPC/E<sup>40</sup> is a widely used model showing excellent predictions for a number of water properties in the liquid region.<sup>41</sup> However, its bad a performance in locating the temperature

of maximum density (TMD), the melting temperature,  $T_m$ , and the isothermal compressibility minimum<sup>42</sup> seems to discourage its use to investigate the supercooled region. Since TIP5P<sup>43,44</sup> provides very good estimates of both the TMD and  $T_m$ , it has been used in simulation studies attempting to disclose the behavior of metastable liquid water.<sup>8,28,36,45,46</sup> However, the excellent performance of the model at ambient conditions is not preserved when one moves away from this region. This failure is particularly serious because it is a signal that the results for the response functions cannot be satisfactory.<sup>41</sup> Moreover, TIP5P gives a very poor estimate of the density of hexagonal ice and, hence, of the density and other properties of a possible low density phase in a liquid-liquid coexistence.

It has been demonstrated<sup>47–49</sup> that the TIP4P geometry is more appropriate than that of three-site models—such as SPC—or five-site models—like TIP5P—to account for the TMD and the liquid-solid equilibrium of water. These studies followed an increasing interest in re-parametrized TIP4P models.<sup>50–52</sup> Among these, TIP4P/2005<sup>52</sup> seems to produce a better overall agreement with the experiment for a large number of properties of water in condensed states.<sup>41,53</sup> Moreover, TIP4P/2005 results are quite accurate for properties relevant to the study of metastable water, namely, water anomalies<sup>42</sup> and the equation of state of supercooled water.<sup>54</sup> Finally, the model gives a quantitative account of recent measurements of the speed of sound of doubly metastable (supercooled and stretched) water.<sup>20</sup>

From the above arguments, it seems then that TIP4P/2005 is the ideal candidate for the study of the water anomalies in the supercooled region. It may come as a surprise that only a reduced number of works have been devoted to this issue,<sup>36,55–62</sup> probably because the model has also some limitations mainly derived from the large structural relaxation times at deeply supercooled states. Abascal and Vega<sup>55</sup> proposed that the model exhibits a LLCP at 193 K and reported a case of a liquid-liquid separation (low- and high-density) below the second critical point. Even though Overduin and Patey<sup>58</sup> argued that longer simulations (8  $\mu$ s for 500 molecules, instead of 400 ns) were necessary to obtain well converged density distributions at those conditions, a number of authors<sup>36,56,59,60</sup> confirmed the rest of the results presented in Ref. 55. The study of Overduin and Patey does not essentially contradict the results of Abascal and Vega if the suggested LLCP of TIP4P/2005 would be slightly shifted towards lower temperatures. This is in line with the critical temperature reported for this model by Sumi and Sekino<sup>56</sup> (182 K) and Yagasaki *et al.*<sup>36</sup> (185 K). Interestingly, a two-structure equation of state consistent with the presence of a LLCP provides a very similar critical temperature.<sup>59,62</sup> Therefore, it would be of great interest to perform a simulation study using advanced sampling methods to unambiguously check the existence of a LLCP for this model (similar to that successfully accomplished for ST2<sup>35</sup>). However, the work of Overduin and Patey clearly indicates that such study would be extremely costly in computer time.

In this work, we propose to circumvent the question of the existence of a LLCP and focus on the related issue of the shape of the TMD and its relation to other water

anomalies. As shown above, the study not only involves calculations in the supercooled and/or stretched region *outside* the proposed critical region but it may also provide a complete perspective of the scenario of water anomalies. Although the study is highly demanding in computer resources, it is still affordable. We also note that regardless of the nature of the phase diagram of TIP4P/2005 liquid at low temperature (real or virtual critical point,...), the region we consider in the present paper includes the one relevant for experiments on bulk water. It can therefore serve as a guide for future measurements.

## II. METHODS

All simulations (except those intended for the calculation of the vapor-liquid spinodal) have been performed with 4000 TIP4P/2005 water molecules in the isothermal-isobaric  $NpT$  ensemble using the molecular dynamics package GROMACS 4.6<sup>63,64</sup> with a 2 fs time step. Long range electrostatic interactions have been evaluated with the smooth particle mesh Ewald method.<sup>65</sup> The geometry of the water molecules has been enforced using *ad hoc* constraints, in particular, the LINCS algorithm.<sup>66,67</sup> To keep the temperature and pressure constant, the Nosé-Hoover thermostat<sup>68,69</sup> and an isotropic Parrinello-Rahman barostat have been applied<sup>70</sup> with 2 ps relaxation times.

Most of our calculations intended to evaluate the extrema of thermodynamic properties, and we have adapted our strategy according to this goal. First, we calculated the desired property at regular intervals along isotherms/isobars to provide a rough estimate of the position of the maximum or minimum. Then we ran additional points to precisely locate it. We monitored the uncertainties along the simulation and extended the runs until the differences between the consecutive points were larger than the statistical uncertainty. Thus, the required simulation times varied widely for the different properties and state points. Since most of the calculations correspond to regions where the relaxation of the system is quite slow, the length of the simulations is often of the order of a few hundreds of ns, reaching 1.3  $\mu$ s for the longest run. Despite the careful monitoring of the runs to save computer resources, the required simulation times together with the use of a relatively large system size implicate an important computational effort (equivalent to more than 300 000 h of 2.6 GHz Xeon cores) that has been achieved by means of a GPU-based supercomputer.

The uncertainty on each measurement has been calculated using a method proposed by Hess.<sup>71</sup> The trajectory is divided in blocks, the average for each block is calculated, and the error is estimated as the standard deviation of the block averages. Also, an analytical block average curve is obtained by fitting the autocorrelation between block averages to a sum of two exponentials. In this way, the calculated uncertainties lead to an asymptotic curve only if the trajectory is long enough so that the blocks are uncorrelated. In summary, the procedure not only provides an estimate of the error but also sheds light on the convergence of the trajectory. An example of the application of the method is given as supplementary material.<sup>72</sup>

## III. RESULTS

Although TIP4P/2005 provides quite acceptable results for the density of water at positive pressures (also including the supercooled region<sup>54</sup>), its performance at negative pressures has not been thoroughly assessed (see however Refs. 20 and 25). Very recently, an experimental equation of state for water down to  $-120$  MPa has been reported.<sup>25</sup> This allows to check for the first time the predictions for the equation of state in the large negative pressures region. The numerical values of the density along some isobars for the TIP4P/2005 model are given as supplementary material.<sup>72</sup> Figure 1 shows that the agreement between simulation results and experiment is excellent although the departures increase with decreasing pressures. As a consequence, the prediction for the TMD is slightly shifted, the difference at  $-80$  MPa being about  $7^\circ$ .

In the positive pressures region, the simulation data show a maximum density for isobars up to a pressure of 200 MPa. In accordance with the experiment,<sup>73</sup> at increasing pressures, the TMD shifts to lower temperatures. At negative pressures, the slope of the TMD becomes increasingly negative until the curve retraces (point A in Fig. 2). From the turning point to highly negative pressures, the TMD has a positive slope. This result has already been reported in previous works for much smaller samples.<sup>25,74</sup> Our results for the larger system are very similar to the previous ones and indicate that finite size effects in this region (if any) are quite small (see supplementary material<sup>72</sup>). The largest (negative) pressure for which we have been able to calculate the temperature of maximum density is  $-170$  MPa. Unfortunately, at  $-200$  MPa the system cavitated for several runs using 4000 water molecules. However, using 500 molecules allowed us to perform short runs before the system cavitated so an approximate calculation of the densities and the approximate position of the TMD at this pressure is possible.

The isobars at negative pressures also exhibit a temperature of minimum density,  $T_{mD}$ . However, in the positive pressures region, only the 0.1 MPa curve shows the density minimum. In all cases, the minimum is quite shallow and it is barely appreciable, especially in the case of the  $-125$  MPa and  $-170$  MPa isobars. Fig. 3 shows a detail of

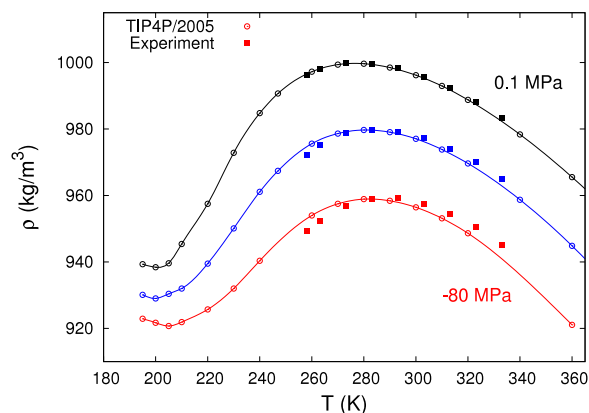


FIG. 1. Densities predicted by the TIP4P/2005 model compared to recent experimental data<sup>25</sup> for the 0.1 MPa,  $-40$  MPa, and  $-80$  MPa isobars.

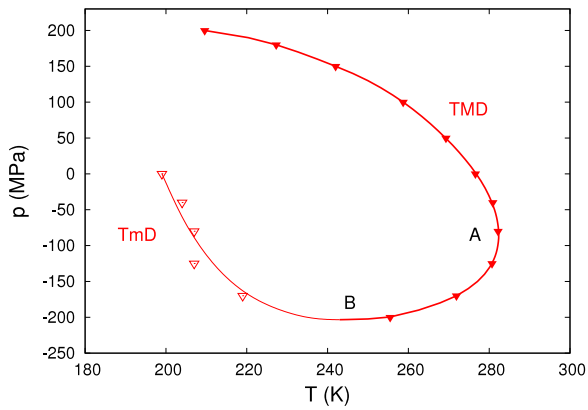


FIG. 2. Locus of density maxima (TMD, thick line) and minima (TmD, thin line). Points A and B mark the turning point of the TMD curve and the point at which the TMD and TmD lines meet, respectively.

these isobars clearly demonstrating the existence of density minima.

Despite the great computational effort and the notable accuracy of the density calculations, the uncertainty of the temperatures of minimum density is about  $5^\circ$  and the loci of the TmD produce a less smooth curve than that of the TMD ones (see Fig. 2).

The existence of a density minimum in real water has not been described previously. Liu *et al.*<sup>75</sup> have reported a density minimum in deeply supercooled confined deuterated water. At ambient pressure, the minimum density occurs at 210 K with a value of  $1041 \text{ kg/m}^3$ . Despite that it is difficult to know how the confinement affects the water properties, we may use these data as a rough guide of the behavior of bulk water. As to TIP4P/2005, previous simulation results indicated the existence of the density minimum<sup>42,56,60</sup> but the accuracy of the data did not allow for a trustworthy estimation of its value. The result of the present work at 0.1 MPa is  $\rho = 938.1 \text{ kg/m}^3$  and is located at  $200 \pm 5 \text{ K}$ . Assuming that the density of deuterated water is 1.106 times that of normal water,<sup>75</sup> we get  $\rho = 1038 \text{ kg/m}^3$ , close to the experimental result in confined water.

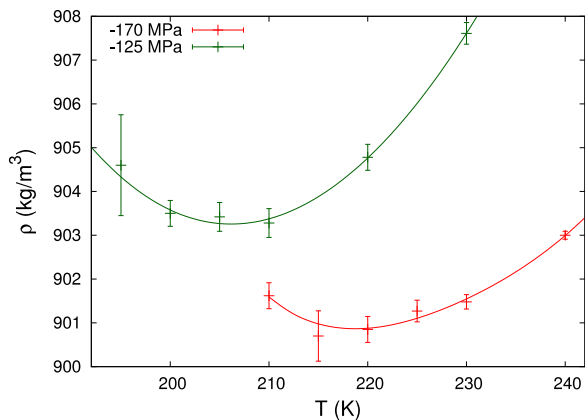


FIG. 3. Detail of the densities for the  $-125 \text{ MPa}$  (decreased by  $0.8 \text{ Kg/m}^3$ ) and  $-170 \text{ MPa}$  isobars showing the existence of a density minimum. The error bars correspond to the 90% confidence interval calculated from the standard deviations obtained in the block procedure. Lines are a weighted fit of the data to a fourth order polynomial.

The difference between the temperature of maximum and minimum density has a peculiar behavior because it is larger near the retracing point of the TMD and decreases at both higher and lower pressures. At large negative pressures, the TMD and TmD lines converge asymptotically (point B in the bottom panel of Fig. 2). Below this pressure, the density no longer exhibits maxima nor minima.

As commented in the introduction, at the retracing point, the TMD curve must be crossed by the line joining the locus of isothermal compressibility extrema. The upper panel of Figure 4 shows  $\kappa_T$  as a function of temperature for several isobars from  $-170 \text{ MPa}$  to  $100 \text{ MPa}$  (for clarity only a few of the simulated isobars are depicted). For pressures higher than about  $-80 \text{ MPa}$ , the isobars show clearly the presence of a maximum and a minimum. At this pressure, the curve exhibits almost imperceptible extrema but, for a slightly lower pressure, the maximum and minimum of  $\kappa_T$  collapse into an inflection point. Then, at large negative pressures,  $\kappa_T$  is a monotonously increasing function of temperature. The locus of  $\kappa_T$  extrema is plotted, in the  $p$ - $T$  plane, in the lower panel of Fig. 4 together with the TMD curve. As expected, both lines cross at the turning point of the TMD (point A in Fig. 4). Notice that the intersection point A lies near the pressure at which  $\kappa_T$  becomes a monotonous function.

On the other hand, theoretical considerations<sup>76</sup> indicate that the locus of isobaric heat capacity extrema along isotherms separates the TMD from the line of minimum densities, TmD. The simulation results for  $C_p$  along isotherms are

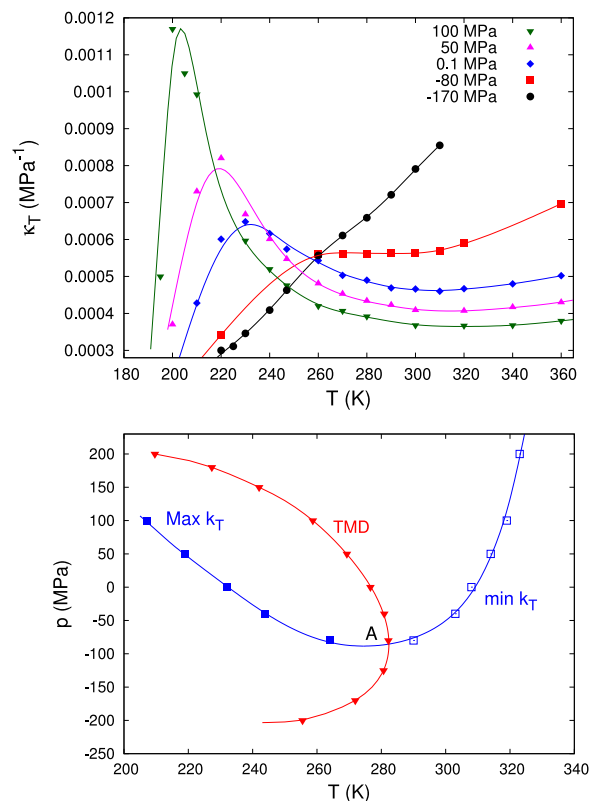


FIG. 4. Top: Isothermal compressibility as a function of temperature along isobars. Bottom: locus of  $\kappa_T$  extrema together with TMD line. Both curves intersect at the point A, the retracing point of the TMD. The curves are cubic splines to guide the eye.

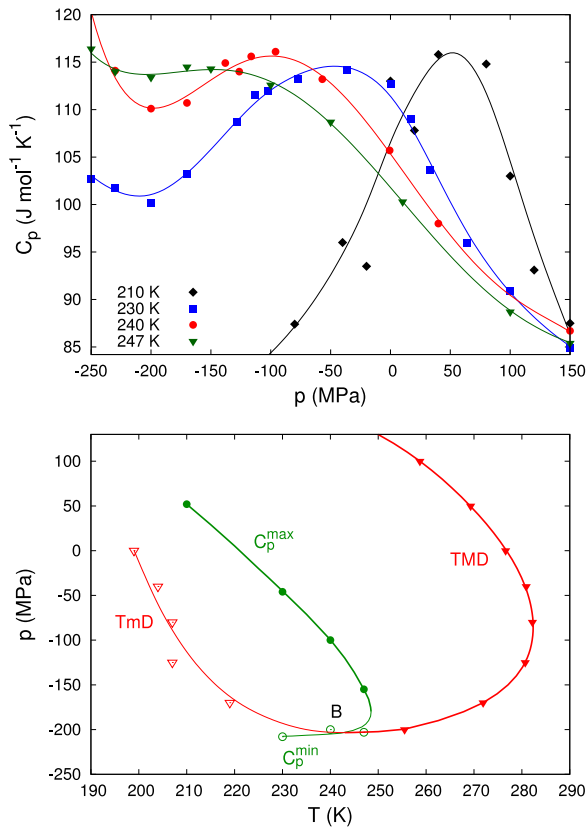


FIG. 5. Top: Heat capacity at constant pressure along isotherms. Bottom: Locus of  $C_p$  extrema together with the curve of density extrema. Both lines intersect at point B separating the TMD and TmD. The curves are a guide to the eye.

presented in Figure 5 (upper panel). The isotherms show both a maximum and a minimum. The separation between these extrema becomes increasingly smaller as the temperature increases. Eventually, at a temperature slightly above 247 K, the maximum and minimum converge to an inflection point and  $C_p$  becomes a monotonous function. The pressures at which the  $C_p$  extrema occur for each isotherm are shown in the lower panel of Fig. 5. In this figure, we have also depicted the locus of density extrema along isobars. As shown by Poole *et al.*,<sup>76</sup> the latter curve must have a zero slope at the intersection point, a condition which is satisfactorily fulfilled by our simulations (see point B in Fig. 5). The location of this point for TIP4P/2005 is about (243 K, -203.4 MPa).

It is clear at this point that the SLC conjecture is not valid for TIP4P/2005 and that the vapor-liquid (VL) spinodal should not meet the (retracing) TMD line. In the SF and LLC scenarios, both curves do not intersect. It is then interesting to check whether this is fulfilled by our calculations. We have tried to evaluate the VL spinodal by locating the zero slope of the pressure-volume curves along isotherms. The simulations were performed in the canonical (NVT) ensemble. However, the system with 4000 molecules sometimes cavitates before providing statistically significant results. We were then forced to reduce the size of the system to 500 molecules. For these samples, the probability of a cavitation event is almost one order of magnitude smaller and it is then possible to obtain statistically significant results. It is well known that finite size

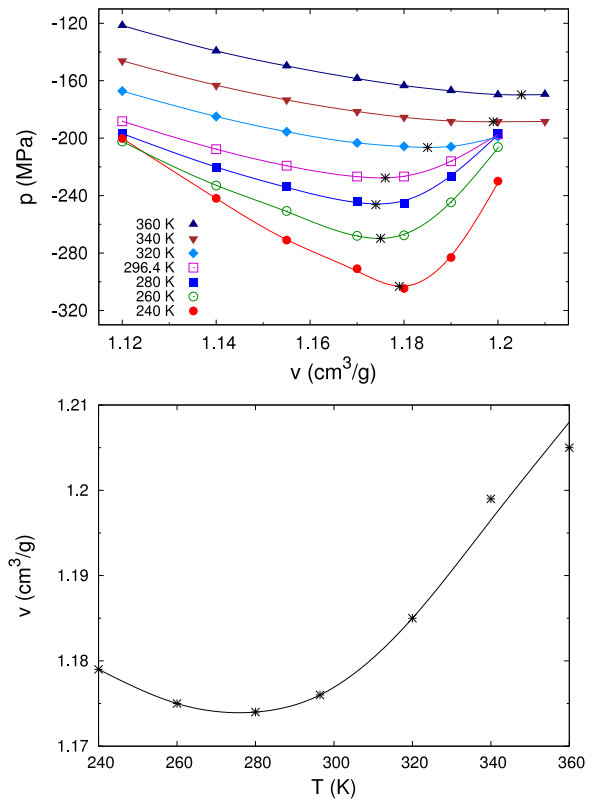


FIG. 6. Estimation of the vapor-liquid spinodal. Top: Pressure as a function of the specific volume for several isotherms (asterisks indicate the position of the minimum). Bottom: Specific volume as a function of temperature along the vapor-liquid spinodal.

effects may be important in this region<sup>77</sup> so our calculations must be seen as a first approximation to the actual spinodal. The corresponding pressure-specific volume isotherms are presented in Figure 6 (top panel) from which we may extract the p-v-T values of the VL spinodal. It is to be noticed that the specific volume along the spinodal shows a non-monotonic dependence on both temperature and pressure (bottom panel of Fig. 6). As expected, the slope of the VL spinodal in the p-T plane is positive and do not meet the TMD curve (Figure 7).

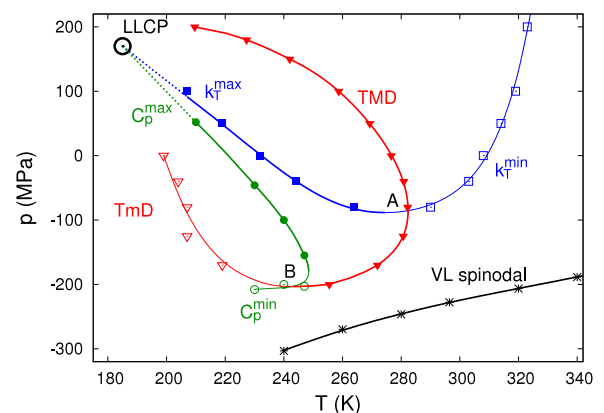


FIG. 7. Comprehensive scenario relating thermodynamic water anomalies and the vapor-liquid spinodal. Dashed lines are an extension of the loci of  $\kappa_T$  and  $C_p$  maxima up to the proposed location of the LLC. <sup>36,62</sup>

The results shown in the bottom panels of Figs. 4 and 5 together with the results of Fig. 6 allow us to give a comprehensive picture of the water anomalies and their relation to the vapor-liquid spinodal. The corresponding plot is presented in Figure 7. Notice that the lines of  $\kappa_T$  and  $C_p$  maxima approach one to another at high pressures and move away as the pressure decreases. In the SF scenario, both curves would only converge at 0 K. Thus, although both the SF and LLC scenarios are compatible with our results, the rate of convergence of these lines seem to favor the LLC hypothesis. This conclusion is also supported by a comparison of Fig. 7 with the corresponding one for ST2.<sup>76</sup> The scenarios of the ST2 and TIP4P/2005 water models are completely analogous and suggest that a critical point for TIP4P/2005 is very plausible. In fact, we have extended the lines of  $\kappa_T$  and  $C_p$  maxima up to the location of the LLC proposed in recent papers.<sup>36,62</sup> The extended lines provide a smooth transition from the critical point to our simulation results for the response functions maxima.

#### IV. CONCLUDING REMARKS

In this work, we have calculated the loci of the extrema of several thermodynamic functions for the TIP4P/2005 water model. In particular, the maxima and minima of the isothermal compressibility along isobars and isobaric heat capacity along isotherms have been evaluated and put in connection with the maxima and minima of the density along isobars. Our work provides for the first time a comprehensive picture of the thermodynamic water anomalies for TIP4P/2005 and their relation to the vapor-liquid spinodal.

The interpretation of previous simulations for the TIP4P/2005 water model in the supercooled region has been rather controversial.<sup>36,55,57,58,61,78</sup> The debate has been mainly focused on three issues: spontaneous phase separation, finite site effects, and ice coarsening. Most of the calculations of this work correspond to the supercooled and/or the stretched regions though we have deliberately avoided the vicinity of the conjectured liquid-liquid region. Our results are then beyond the current debate on the possibility of observing spontaneous liquid-liquid phase separation.<sup>36,61,78,79</sup> This has allowed us to get converged results for the properties of interest with a large but affordable computational effort.

The size of the system, 4000 water molecules, seems to ensure that our calculations are free of finite size effects. It has been reported that even larger samples could be needed for temperatures below the proposed LLC.<sup>61</sup> However, our calculations for the TMD indicate that the differences between the results obtained with 500 and 4000 molecules are marginal (see Fig. 4 of the supplementary material<sup>72</sup>). Thus, at least at the thermodynamic conditions of this work, we do not observe a significant system size dependence.

It has been argued that the phenomenon suggesting the metastability of two distinct liquid phases is actually the coarsening of the ordered ice-like phase.<sup>57</sup> Again, the range of temperatures and pressures of this work indicate that our results are free of the problem of ice coarsening. In a recent study, Espinosa *et al.*<sup>80</sup> have calculated the size of the critical cluster and the nucleation rate for the crystallization

of TIP4P/2005 water as a function of the supercooling. The results for both magnitudes indicate beyond any doubt that our simulations correspond to a metastable liquid. Inherent to metastability is the formation and breaking of small clusters of the stable phase. Thus the question is not the appearance of small crystal nuclei but whether a critical cluster may appear in the simulation. Espinosa *et al.* have evaluated the supercooling required for the formation of a single critical cluster in a simulation with a box side of 40 Å (corresponding to a typical supercooled water density of about 0.94 g/cm<sup>3</sup> in a system of 2000 molecules) for 1 μs. At these conditions (very similar to those of our longest simulations), they report a 65 K supercooling. Since the melting temperature of the model is around 250 K, the appearance of a critical cluster above 185 K is a very unlikely event (notice that the lowest temperature of our calculations is 195 K).

Although in this work we have avoided the issue of the existence of a LLC, it is evident that the overall picture is consistent with both the SF and LLC conjectures. However, the way in which the lines of maximum  $\kappa_T$  and  $C_p$  approach one to another seems to indicate that they meet not too far from the region of calculations clearly favoring the LLC hypothesis over the SF one. This idea is reinforced when one observes that the scenario presented in Fig. 7 very much resembles that of ST2<sup>76</sup> for which most authors give for demonstrating the existence of a LLC.

It is important to stress that the significance of this work goes beyond the theoretical interpretation of simulation results. Most of the thermodynamic states relevant to this work correspond to the negative pressures region where the water properties are largely unknown.<sup>81</sup> Since the TIP4P/2005 water model has demonstrated to provide semiquantitative predictions of the water properties in the supercooled and/or stretched regions,<sup>20,25,54</sup> the scenario of the water anomalies predicted by this model may provide a hint as to where the long-sought for extrema in response functions might become accessible to experiments.

*Note added in proof.* After sending the accepted version of this work we have been aware of a paper by Lu *et al.*<sup>82</sup> reporting a similar study using the coarse grained mW and mTIP4P/2005 water models.

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